

# **EXHIBIT 15**

# MICROELECTRONIC DEVICES

**2nd Edition**

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### *Operating regions of the characteristics*

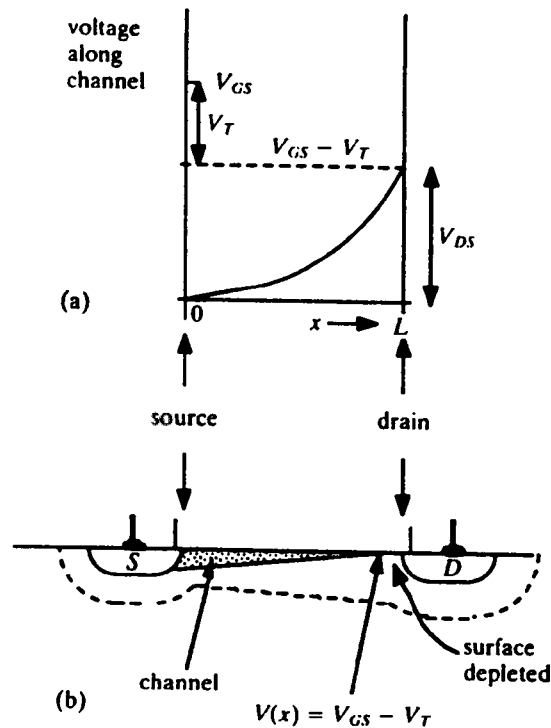
The region of Fig. 4.7(a) to the left of the dashed line, in which eqn. (4.4) remains valid, is often known as the TRIODE REGION of the characteristics. In this book we shall refer to it for clarity as the UNSATURATED REGION, to distinguish it clearly from the SATURATION REGION, where the drain current does not increase with  $V_{DS}$ . The initial portion of the the unsaturated characteristic, where  $V_{DS} \ll V_{GS} - V_T$ , is sometimes called the LINEAR region of operation.

A third principal region of the characteristics, in which the transistor is OFF, (i.e. no drain current flows), exists for all gate voltages below threshold. This can be seen in Fig. 4.7(b), which shows how the drain current  $I_D$  changes with gate-source voltage  $V_{GS}$  when  $V_{DS}$  is held constant at a small value. One practical point to note about this diagram is that the threshold of conduction is not as sharp as the above discussion would imply. The threshold voltage is normally defined for a real device by extrapolating the linear part of this characteristic back to the voltage axis. The reason for the curvature near this point is covered in Panel 4.2, which should be read following the discussion of the saturation region of the characteristics in the next section.

#### **4.4 The saturation region of the output characteristic**

When  $V_{DS} \geq V_{GS} - V_T$  the transistor is said to be SATURATED, meaning that the drain current no longer rises strongly with  $V_{DS}$ . At the onset of saturation, the potential relative to the source varies along the channel, as shown in Fig. 4.8(a), reaching the value  $(V_{GS} - V_T)$  at the drain. At this point, the voltage between gate and drain,  $V_{GD}$ , just equals the threshold voltage  $V_T$ . The channel is said to be PINCHED OFF at the drain, and the drain-source voltage  $V_{DS}$  at which this occurs is often referred to as the pinch-off voltage. Note, however, that the pinch-off voltage does not have a fixed value, as it is equal to  $(V_{GS} - V_T)$ .

Any further increase in  $V_{DS}$  ensures that the channel becomes fully depleted of carriers at the end nearest the drain. Thus, an extremely short depletion region forms between the channel itself and the drain, as illustrated in Fig. 4.8(b). The electrons in the channel must flow across this depleted region in order to reach the drain. They can do this with the help of even



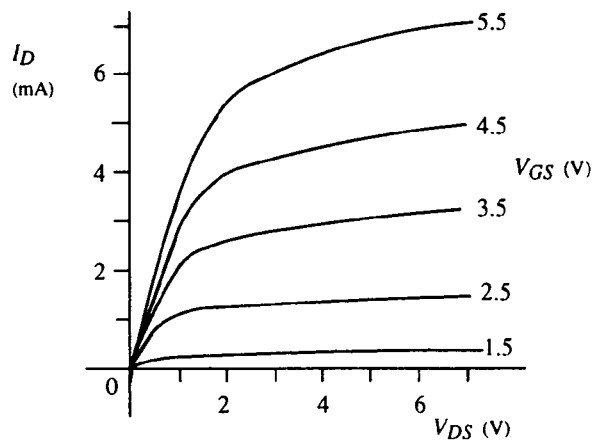
**Fig. 4.8** (a) Voltage variation along the channel in saturation (b) a depletion region of length  $L_D$  is formed in saturation at the drain end of the channel.

a small voltage difference, because the distance involved is so small that a large electric field strength is easily created. At the point where the conducting channel ends (is 'pinched-off'), the potential is equal to the pinch-off voltage ( $V_{GS} - V_T$ ). Note that the difference between the voltage  $V_{DS}$  and the pinch-off value of ( $V_{GS} - V_T$ ) appears across the short, high-resistance depletion layer, and is enough to give the few electrons present there a high velocity towards the drain. Indeed, they travel at their maximum speed, the *saturation velocity*.

We expect the current to rise no further, because an increase in  $V_{DS}$  no longer raises the velocity of carriers in the conducting channel itself, but merely lengthens the depleted region slightly. The drain current thus saturates, and the characteristics are shown horizontal on the theoretical plots in Fig. 4.7(a). The drain current in saturation is simply found by putting  $V_{DS} = (V_{GS} - V_T)$  into eqn. (4.4), thus

$$I_D = \frac{1}{2} \mu_e C_{ox} \frac{W}{L} (V_{GS} - V_T)^2 \quad \text{in saturation} \quad (4.5)$$

However, constancy of the saturation current depends upon maintaining a constant channel length. In practice, the depletion layer at the drain end lengthens as the voltage across it increases. The channel length  $L$  decreases slightly in consequence, and the drain current rises a little. This is known as the EARLY effect, after the first person to describe the rather similar behaviour observed in bipolar transistors. It is also referred to as CHANNEL LENGTH MODULATION, and it will be explained at greater length in section 4.6. Figure 4.9 shows some transistor characteristics where the effect is clearly seen.



**Fig. 4.9** Transistor characteristics showing the effect of channel length modulation in the saturation region.

We can use eqn. (4.5) to find the boundary between the unsaturated and saturation regions on Fig. 4.7(a) by naming the values of  $I_D$  and  $V_{DS}$  at the onset of saturation as  $I_{DSat}$  and  $V_{DSat}$  respectively. Then, since  $V_{DSat} = V_{GS} - V_T$ , eqn. (4.5) gives

$$I_{DSat} = \frac{1}{2} \mu_e C_{ox} \frac{W}{L} V_{DSat}^2$$

This is the equation of a parabola, which is shown dashed in Fig. 4.7(a).

It is worth noting that the term saturation has a quite different meaning in the case of the bipolar transistor, as will be explained in chapter 5.

## PANEL 4.2

*The pinched-off channel: subthreshold conduction*

When  $V_{DS} \geq V_{GS} - V_T$ , i.e. when the *drain* end of the channel is pinched off, the carrier concentration there is not zero but merely small in magnitude, just as in the depletion region of a  $p-n$  junction. The gate voltage is unable to maintain *strong inversion* at the drain end of the channel; we say that it supports *weak inversion*. Similarly, weak inversion occurs at the source if we make  $V_{GS}$  slightly less than  $V_T$ : the transistor is off and, in the normal way, we would then say that the transistor is *below threshold*.

The term 'weak inversion' implies that the channel region *does* contain some free electrons. This is the so-called sub-threshold region of operation, in which the electron concentration in the channel is sufficient only to allow very small currents to flow. In this regime, the following approximate expression for the drain current  $I_D$  applies, provided that  $V_{SB}$  and  $V_{DB}$  are not too close to the pinch-off voltage  $V_{GS} - V_T$ , which is abbreviated to the symbol  $V_p$ :

$$I_D = 2\mu C_{ox} \frac{W}{L} \left( \frac{kT}{e} \right)^2 \left\{ \exp\left( \frac{V_p - V_{SB}}{kT/e} \right) - \exp\left( \frac{V_p - V_{DB}}{kT/e} \right) \right\}$$

Here the source and drain voltages are expressed relative to the substrate, or bulk silicon.

More accurate (but still relatively simple) equations for  $I_D$  are available for use when  $V_{SB}$  and  $V_{DB}$  are close to  $V_p$ .

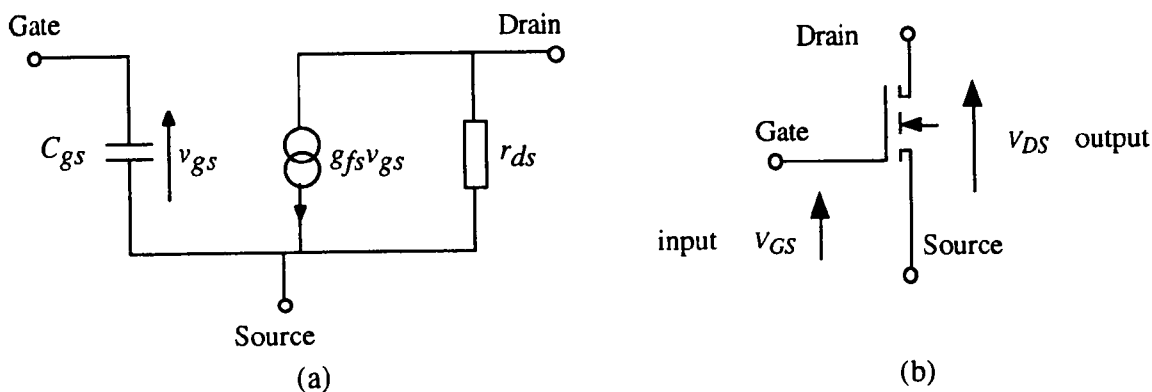
The sub-threshold region is important commercially for two reasons:

- (i) it determines the 'off' current of a transistor in many digital circuits, and hence the quiescent current drawn by each of the millions of switches used in a complete CMOS integrated circuit, and
- (ii) it can be used instead of the unsaturated region for low-power applications (e.g. in watches, satellites and body implants), and is likely to become more important due to a growing appreciation by more engineers of the circuit design techniques needed for its use.

## 4.5 The small-signal equivalent circuit

As with any electronic device, an equivalent circuit can be constructed to model the way small changes in input voltage affect both the input current and the output voltages and currents. It is used to represent the behaviour of the transistor when it is employed in a circuit designed to amplify small currents or voltages at the input terminals. The simplest equivalent circuit for a MOSFET is shown in Fig. 4.10(a), alongside the symbol of the transistor which it models. Note the arrow pointing from the  $p$ -type substrate to the  $n$ -channel, which identifies the channel type.

The convention followed in chapter 3 for the use of small letters with small subscripts for small-signal voltages and currents is continued here. The circuit of Fig. 4.10(a) is used to model the transistor in the connection shown in Fig. 4.10(b). This is known as the common-source connection, since the source terminal is common to both the input and output.



**Fig. 4.10** (a) Small-signal equivalent circuit of a MOSFET in the common-source connection  
(b) Circuit symbol of the  $n$ -channel MOSFET, shown in the common-source connection.

It is necessary to treat the unsaturated and saturation operating regions of the transistor separately, as in earlier sections. Within each region of operation, the component values depend upon the bias voltages, i.e. upon the dc components of the voltages  $V_{GS}$  and  $V_{DS}$ .

When used in an amplifier, a MOSFET is usually biased in its saturation region, so we shall concentrate more upon the saturation region in what follows.



Let us first consider how to define each of the component values in Fig. 4.10(a).

*The gate-source input capacitance  $C_{gs}$*

A MOSFET has almost infinite input resistance, owing to the excellent insulating properties of the gate oxide. Hence, we normally need only to take account of the input capacitance.

The input capacitance can be defined as the small-signal charge  $q_g$  on the gate electrode divided by the small-signal gate-source voltage  $v_{gs}$ . Thus

$$C_{gs} = \frac{q_g}{v_{gs}} = \frac{\partial Q_G}{\partial V_{GS}}$$

*The transconductance  $g_{fs}$*

The dependence of the small-signal drain current  $i_d$  on small-signal gate-source voltage  $v_{gs}$  is modelled by a current generator of strength  $g_{fs}v_{gs}$  in Fig. 4.10, where  $g_{fs}$  is given by

$$g_{fs} = \frac{i_d}{v_{gs}} = \left( \frac{\partial I_D}{\partial V_{GS}} \right)_{V_{DS} = \text{const}}$$

$g_{fs}$  is referred to as the forward TRANSCONDUCTANCE\* of the transistor in the common-source connection. It has the dimensions of a conductance, but is unlike an ordinary conductance in that it connects an *output* current with an *input* voltage. The two subscripts *f* and *s* are used to distinguish this transconductance from others which are sometimes needed, for example the reverse transconductance  $g_{rs}$ , used when the drain and source are interchanged, and similar transconductances defined for cases when the terminal which is common to both input and output, is not the source (denoted by the second subscript *s* to the symbol), but the gate or drain.

\*The symbol  $g_m$  is often used in place of  $g_{fs}$ , which is the internationally agreed symbol for this component.

Note from the equation that in evaluating the transconductance,  $V_{DS}$  is to be held constant, i.e.  $v_{ds} = 0$ . Otherwise, some of the small-signal current generated would flow in the resistance  $r_{ds}$  in the figure. If we are to measure the strength of the current source in Fig. 4.10 by measuring the external drain current, it must be measured in a short-circuit across the terminals.

### *The output resistance $r_{ds}$*

The resistance  $r_{ds}$  in Fig. 4.10 is commonly known as the OUTPUT RESISTANCE of the transistor, since the drain and source are the output terminals of the amplifier. It models that part of the change in drain-source voltage  $V_{DS}$  which is proportional to  $I_D$ . Thus the definition of  $r_{ds}$  is

$$r_{ds} = \frac{v_{ds}}{i_d} = \left( \frac{\partial V_{DS}}{\partial I_D} \right)_{V_{GS} = \text{const}}$$

We shall use these definitions to find expressions for the component values in the unsaturated region of operation, since this is the easiest case to consider. The saturation region requires a little more development of the ideas introduced in earlier sections.

## **4.6 Component values in the equivalent circuit: the unsaturated region of operation**

### *The gate-source input capacitance*

The two input terminals, the gate and the source, present a very high d.c. input resistance, typically above  $10^{12} \Omega$  and a small capacitance  $C_{gs}$  in the range 0.01–10 pF, which is shown in Fig. 4.10. The input resistance is sufficiently high that for practical purposes it is commonly omitted, as here. Hence at all useful frequencies above d.c., the input impedance is

$$Z_{in} = (j\omega C_{gs})^{-1}$$

Since  $C_{gs}$  is small, the input impedance has a very large magnitude, and indeed may be assumed infinite for practical purposes, except at high frequencies, in all regions of operation. We shall discuss what determines the value of  $C_{gs}$  later, along with other capacitances which play a role at high frequencies.

### *Transconductance in the unsaturated region of operation*

The definition of  $g_{fs}$  above is easily used in conjunction with eqn. (4.4) for  $I_D$  to find an expression for the transconductance in the unsaturated region. Differentiating eqn. (4.4) gives the result

$$g_{fs} = \mu_e C_{ox} \left( \frac{W}{L} \right) V_{DS} \quad (\text{unsaturated})$$

### *Output resistance in the unsaturated region*

By differentiating eqn. (4.4) with respect to  $V_{DS}$  it is easy to show that

$$r_{ds} = \mu_e C_{ox} \left( \frac{W}{L} \right) [V_{GS} - V_T - V_{DS}] \quad (\text{unsaturated})$$

## **4.7 Component values in the equivalent circuit: the saturation region of operation**

### *Input capacitance*

The d.c. input resistance remains effectively infinite.

The input impedance is identical to that given for the unsaturated region, thus:

$$Z_{in} = (j\omega C_{gs})^{-1}$$

although the value of  $C_{gs}$  differs somewhat from that in the unsaturated region, as will be discussed later.

### Transconductance in the saturation region

To evaluate the transconductance in the MOSFET's saturation region, we differentiate eqn. (4.5) with respect to  $V_{GS}$ , and find that

$$g_{fs} = \mu_e C_{ox} \frac{W}{L} (V_{GS} - V_T) \quad (\text{in saturation}) \quad (4.6)$$

Here,  $V_{GS}$  denotes the mean gate-source voltage difference (i.e. the d.c. or bias component). Unless the input voltage  $v_{gs}$  is small compared to  $V_{GS}$ , the transconductance  $g_{fs}$  varies with signal amplitude, thereby creating a distortion of the collector current waveform. Note that the designer can vary  $g_{fs}$  at will by choosing the ratio  $W/L$ .

To estimate a typical value of  $g_{fs}$ , we first estimate  $C_{ox}$ . The relative permittivity  $\epsilon_{ox}$  of  $\text{SiO}_2$  is close to 4.0. Thus, assuming an oxide thickness  $t_{ox}$  of  $0.05 \mu\text{m}$ , we have

$$C_{ox} = \frac{\epsilon_{ox} \epsilon_0}{t_{ox}} = \frac{4 \times 8.8 \times 10^{-12}}{5 \times 10^{-8}} = 7.0 \times 10^{-4} \text{ pF}/\mu\text{m}^2.$$

The next point to note is that the mobility  $\mu_e$  of electrons is about one half of the value used in Chapter 1 appropriate to the doping concentration in the channel, because of extra collisions of the electrons with the surface of the Si at the interface with the  $\text{SiO}_2$ . The channel thickness is comparable to the mean free path, so that such collisions are relatively frequent. For example, when the doping is low,  $\mu_e$  is reduced to about  $0.07 \text{ m}^2/\text{Vs}$ . Hence when  $W/L = 10$ , we have typically

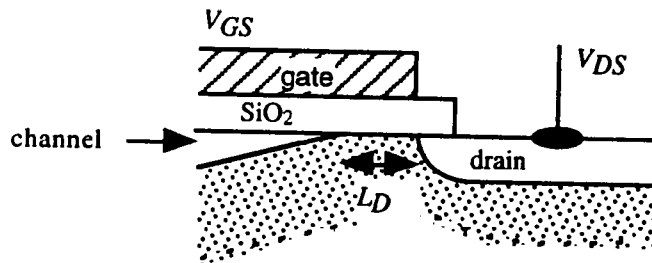
$$\begin{aligned} g_{fs} &= 0.07 \times 7.0 \times 10^{-4} \times 10 (V_{GS} - V_T) \quad (\text{in saturation}) \\ &= 0.28 \text{ mS for every volt by which } V_{GS} \text{ exceeds } V_T. \end{aligned}$$

When compared with the transconductance of a bipolar transistor evaluated in Chapter 5, this is a rather low value when voltages of just a few volts are used.

*Output Resistance in the saturation region (The Early effect, or channel length modulation)*

The output resistance cannot be calculated directly from eqn. (4.5) for the drain current, since  $V_{DS}$  does not appear explicitly there. However, as already explained, the drain current does increase slightly with  $V_{DS}$ .

As we have seen, when the gate-drain voltage  $V_{GD}$  falls below the value of the threshold voltage  $V_T$ , the drain current saturates because the channel at the drain end becomes depleted of carriers. As the drain voltage rises further, with  $V_{GS}$  constant, the point in the channel at which pinch-off begins moves slightly away from the drain, as shown in Fig. 4.11.



**Fig. 4.11** An enlarged cross-sectional view of the channel at the drain end where, in saturation, the channel is pinched-off (i.e. depleted) over a length  $L_D$ .

This reduces the effective channel length of the depleted region by the length  $L_D$ , so that the drain current may be found by modifying eqn. (4.5) as follows

$$I_D = \frac{1}{2} \frac{\mu C_{ox} W}{(L - L_D)} (V_{GS} - V_T)^2 \quad (4.7)$$

As  $L_D$  rises with increasing  $V_{DS}$ ,  $I_D$  also rises slightly, giving the transistor its finite output resistance  $r_{ds}$ . The dependence of  $L_D$  upon  $V_{DS}$  is not easily calculated, however, and approximations are usually used.

The simplest model of this effect, which is provided within the SPICE computer modelling package, assumes that  $L_D = \Lambda V_{DS}$ , where  $\Lambda$  is an adjustable parameter chosen by the user, independent of transistor voltages or currents. The corresponding value of  $r_{ds}$ , the small-signal output resistance

is then simply found by differentiating eqn. (4.7) to give

$$r_{ds} = \left( \frac{\partial I_D}{\partial V_{DS}} \right)^{-1} = \frac{(L - L_D)}{\Lambda I_D} = \frac{V_A}{I_D}$$

where  $V_A = (L - L_D)/\Lambda$  is called the EARLY VOLTAGE. This parameter is approximately constant whenever  $L \gg L_D$ , i.e. in all except transistors with very short channels. Although there is no simple justification for the assumption that  $L_D = \Lambda V_{DS}$ , the simple expression for  $r_{ds}$  which results does indeed describe the behaviour of real transistors fairly well.

Since the voltage across the depleted channel length  $L_D$  is  $V_{DS} - V_p$  (where  $V_p$  is the pinch-off voltage, equal to  $V_{GS} - V_T$ ), then by applying conventional abrupt  $p$ - $n$  junction theory,  $L_D$  might be expected to rise in proportion to  $(V_{DS} - V_p)^{1/2}$ . This is not often the case in practice, since the conventional theory applies only to the doping profile of a planar, abrupt junction, whereas the doping here usually varies with depth as well as distance along the channel. Most versions of SPICE usually provide the option of using this, more complicated, model for the Early effect, in spite of its limitations.

## 4.8 Small-signal capacitances and transient behaviour

Before discussing the small-signal circuit components which model the high-frequency behaviour of a MOSFET, it is helpful to consider the factors which limit the rate at which the drain current responds to a change in the gate-source voltage. The most fundamental of these is the so-called transit time.

### *The transit time*

The maximum rate at which the drain current  $I_D$  can be switched by a change in input voltage is determined by the rate at which the charge in the channel is changed, because it is this charge which carries the current.

Suppose an  $n$ -channel enhancement MOSFET is initially in its unsaturated (triode) region, with  $V_{DS}$  held constant. If  $V_{GS}$  is increased instantly, the